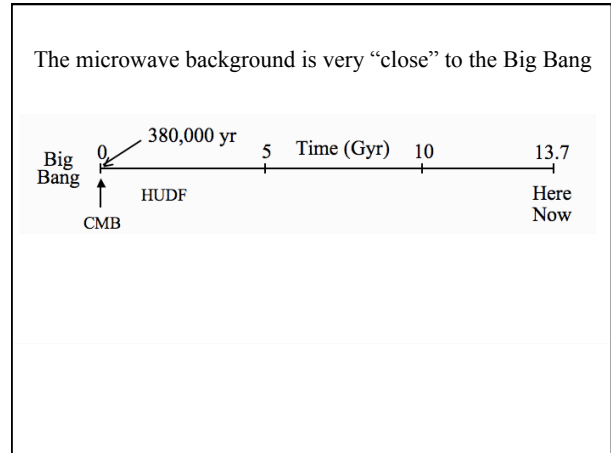
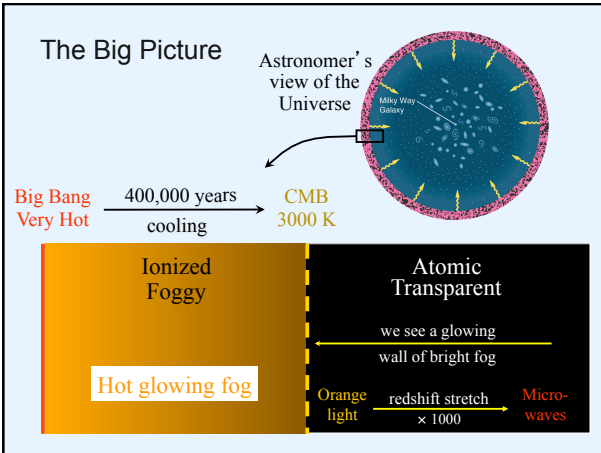
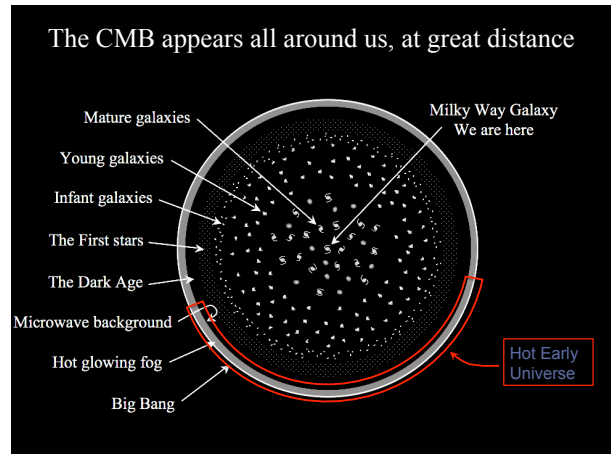


# The Microwave Background

Notes based on Teaching Company lectures, and associated undergraduate text – with some additional material added.

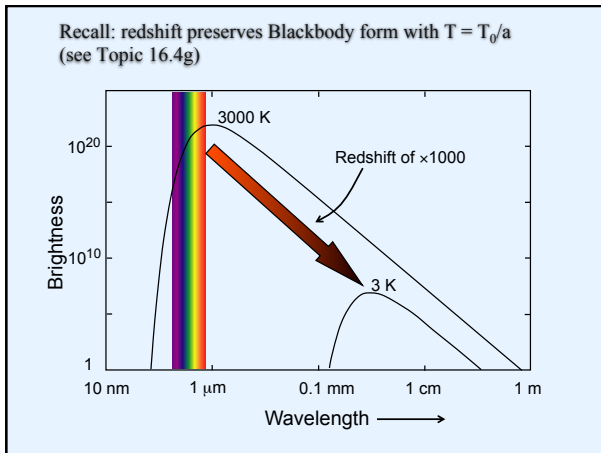
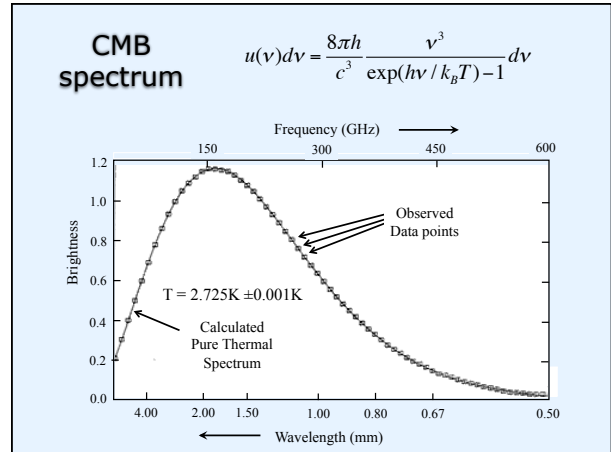
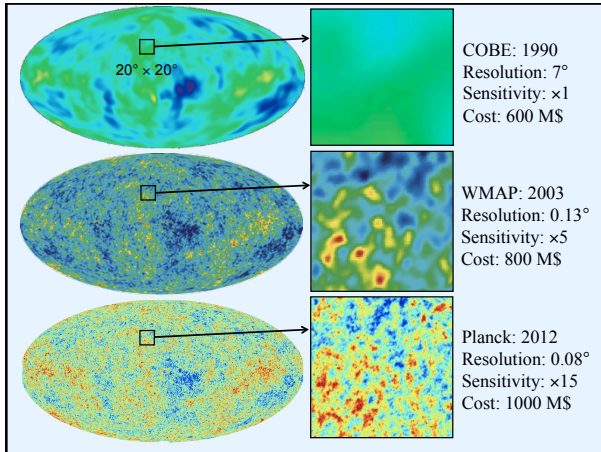


## Discovery of the CMB

In 1964 Penzias and Wilson found an excess of microwave emission coming from the sky. They could not account for its origin.....

Wilson Penzias

## Major CMB Space Observatories



### Temperature of recombination

Naively, recombination might happen when  $k_B T \sim 13.6$  eV (typical photon energy  $\sim$  H binding energy). i.e.  $T \sim 150,000$  K.

However, photons more abundant than protons, so even at lower temperature, there are many more photons with  $E > 13.6$  eV.

Photon/baryon ratio is large and preserved during expansion:

$$\frac{n_\gamma}{n_b} \approx \frac{\Omega_{r,0} / \langle E_\gamma \rangle}{\Omega_{b,0} / \langle E_b \rangle} \approx \frac{5.0 \times 10^{-5} / 7 \times 10^{-4} \text{ eV}}{0.04 / 938 \text{ MeV}} \approx 1.7 \times 10^9$$

At 5700K, the tail of the black body contains as many photons with  $E > 13.6$  eV as there are protons.

### Temperature of recombination

More exactly, we use the Saha equation (for pure hydrogen):

$$\frac{1-X}{X^2} \approx 3.8 \frac{n_b}{n_\gamma} \left( \frac{k_B T}{m_e c^2} \right)^{3/2} \exp\left( \frac{13.6 \text{ eV}}{k_B T} \right)$$

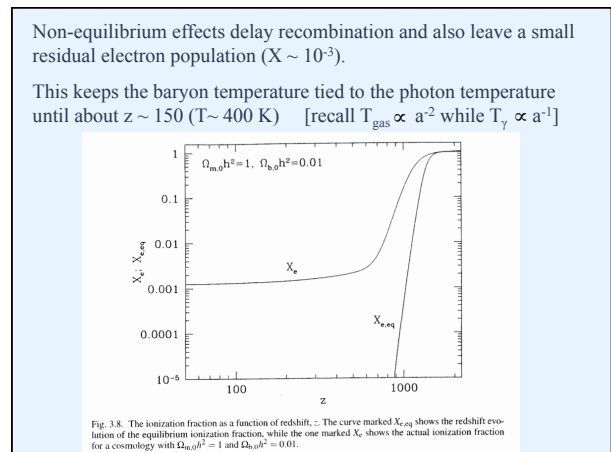
Where  $X$  is the fractional ionization:  $n_e/n_b = n_\gamma/(n_b + n_H)$

Setting  $X \sim 0.5$  we find  $T \sim 3740$  K, or  $z \sim 1370$ , or 240 kyr.  
Transition is quick:  $X = 0.9 - 0.1$  from  $z = 1475 - 1255$  ( $\Delta t \sim 70$  kyr)

Several complicating factors delay this time:

- 1) Most recombination photons are reabsorbed/reionize.
- 2) Reaction rates become longer than expansion rate – reactions fall out of equilibrium.

Ultimately,  $2S \rightarrow 1S$  via 2-photon decay secures recombination.  
Growing m.f.p. allows Ly- $\alpha$  photons to redshift out of resonance.



Note: same physics sets ionization degree in stellar atmospheres. G stars quite neutral; O stars: highly ionized.

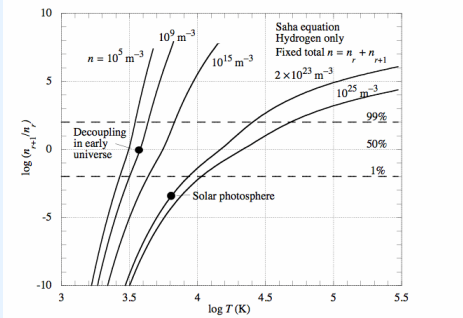


Figure 3. Plot of log ionization fraction  $n_{e+1}/n_e$  for various fixed total densities,  $n = n_e + n_{e+1}$  (i.e.,  $n = n_H + n_p$ ) as a function of temperature.

## Time of decoupling & last scattering

Photons Thomson scatter off free electrons, at a rate:

$$\Gamma_T = n_e \sigma_T c = X_e n_b \sigma_T c$$

When  $\Gamma_T(z) > H(z)$  then the photons no longer interact (decouple). This happens around  $z \sim 1100$  or  $\sim 3000\text{K}$  or  $380,000$  years.

This is very similar to the time of “last scattering”, which is also the surface of optical depth,  $\tau = 1$ , as viewed from  $z = 0$  (here).

$$\tau(z) = \int_0^z \Gamma_T \frac{dt}{dz} dz = \int_0^z \frac{\Gamma_T(z)}{H(z)} \frac{dz}{1+z} \approx 0.37 \left( \frac{z}{1000} \right)^{14.25}$$

This also happens around  $z \sim 1100$ . There is a width  $\Delta z \sim 80$  from which most of the photons last scattered.

This is the “surface” we see as the Microwave Background.

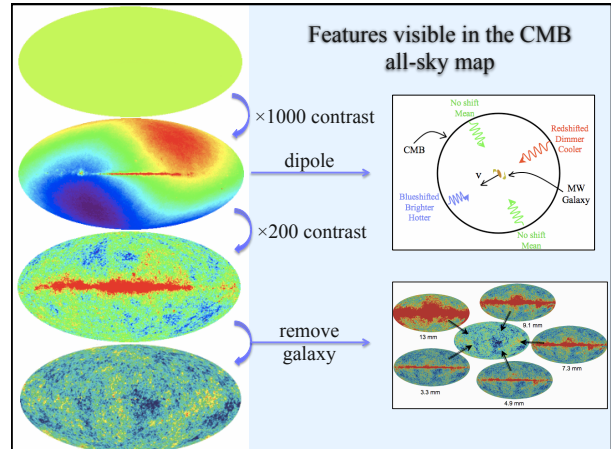
## All-sky CMB image

Since today’s universe is lumpy on small scales, one expects to see variations in the brightness of the CMB.

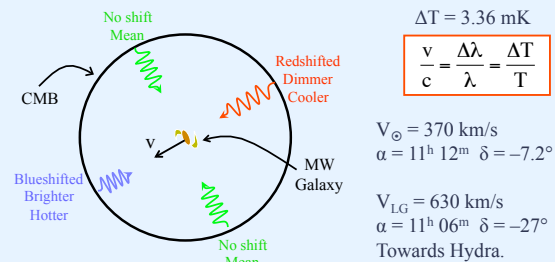
After many years of searching, these were finally seen in the 1990 COBE data.

Since then, mapping these fluctuations has proved to be extremely valuable.

The CMB fluctuations provides our point of entry into the major topic of the origin and growth of structure.

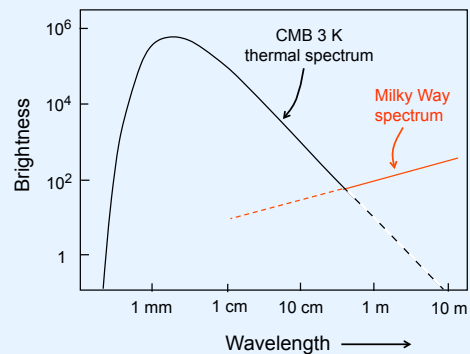


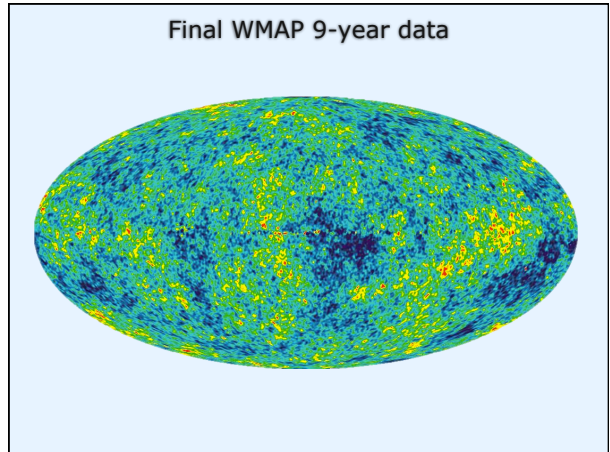
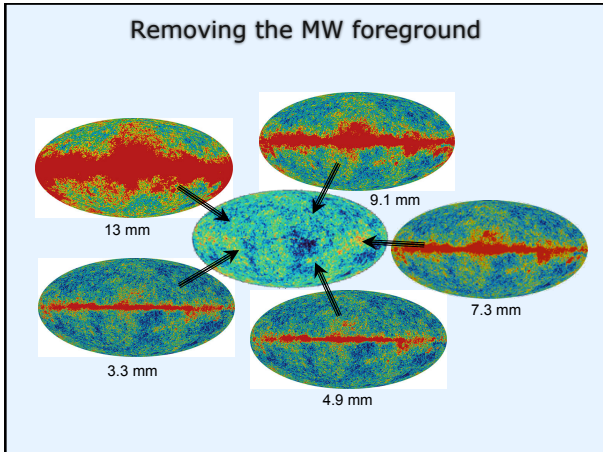
## Doppler origin of the CMB Dipole



The Doppler velocity matches (approximately) an analysis of the local velocity field:  $\sim 300 \text{ km/s}$  to Virgo +  $\sim 500 \text{ km/s}$  to Great Attractor.

## CMB and Galaxy spectra





### Statistical description of anisotropies

Fourier approach to functions on a sphere uses spherical harmonics:

$$\Delta T(\theta, \phi) = \sum_{l,m} a_{l,m} Y_{l,m}(\theta, \phi)$$

Example of  $l = 16$ , summed over  $m$

$$C(l) = \frac{1}{2l+1} \sum_m a_{l,m}^2$$

Normally plot  $C(l) l(l+1)/2\pi \mu K^2$  against  $l \approx 180/\Delta\theta^\circ$ . This shows relative power per log  $l$ .

Note: this function on a surface arises (in part) from a 3-D density field,  $\delta(\mathbf{r}) = \delta\rho(\mathbf{r})/\langle\rho\rangle$ , which is described by  $P(k)$ , with  $k = 2\pi/\lambda$ :

$$\delta(\vec{r}) = \sum_{k_x, k_y, k_z} \delta_k e^{i\vec{k}\cdot\vec{r}}$$

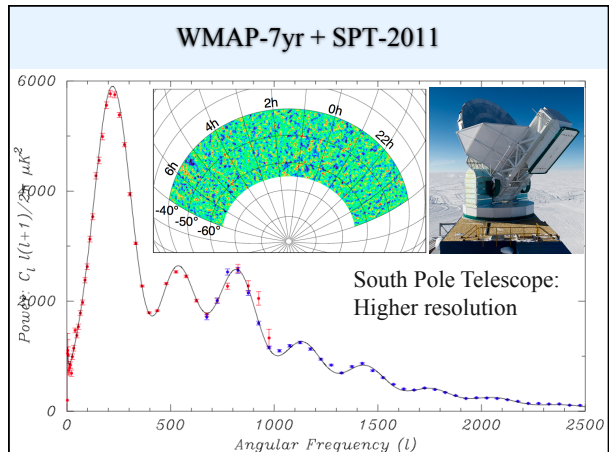
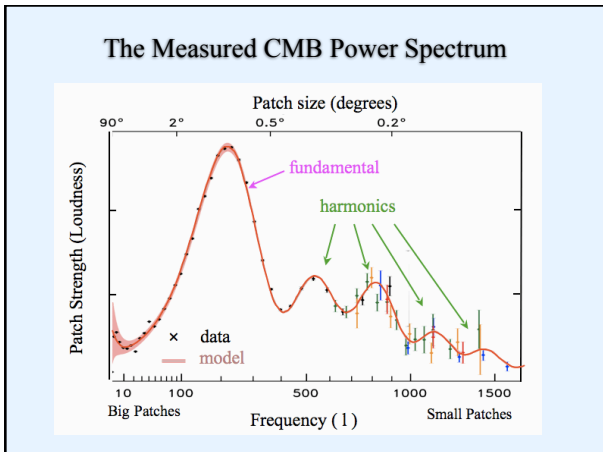
Where  $k = |\mathbf{k}|$  due to isotropy

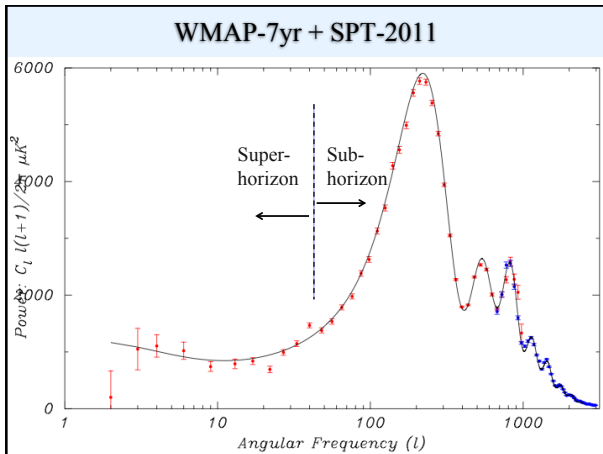
Note: primordial  $P(k) \sim k^{-1}$  has  $C(l) l(l+1) = \text{const}$ .

Note:  $P(k)$  is primordial outside horizon size & modified inside.

### Sky Maps $\Rightarrow$ Sound Spectra

Simple depiction of how the power spectrum is constructed.





### Origin of Anisotropies.

Two main constituents:

- i) dark matter: almost smooth + slight density uneveness
- ii) photon-baryon gas, coupled by Thomson scattering.

Dark matter regions create gravitational valleys.

Gas falls in and bounces out, and falls in again, etc. This is a ~spherical sound wave.

Frequency of sound wave depends on size of region.

### Understanding C(l)

- What are all those bumps and wiggles ??
- Many processes affect C(l), not all are sound waves !!  
Usually divided: **Primary** / **Secondary** / (**Tertiary**)

(1) cold hot (2)  $\Phi$

- **Primary:**
  1. more/less dense (valley/hill)  $\rightarrow$  hot/cold  $\rightarrow$  blue/red
  2. valley/hill  $\rightarrow$  gravitational red/blue shift

these partly cancel:  $\frac{\Delta T}{T} = \frac{\Phi}{c^2}$  &  $\frac{1}{3} \frac{\Delta \rho}{\rho} \Rightarrow \frac{\Delta T}{T} = \frac{2}{3} \frac{\Phi}{c^2}$  Sachs Wolfe

  3. gas falls in / rebounds out  $\rightarrow$  Doppler red/blue shift (sub-horizon only, ie  $l > 50$ )

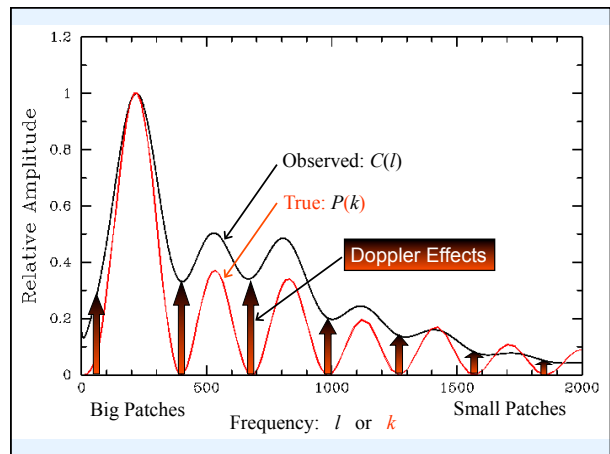
Formation of fundamental & harmonics  $\rightarrow$

### Gravitational Redshift: Sachs-Wolfe Effect

Doppler shifts of the moving gas

blue-shift warmer brighter  
red-shift cooler dimmer

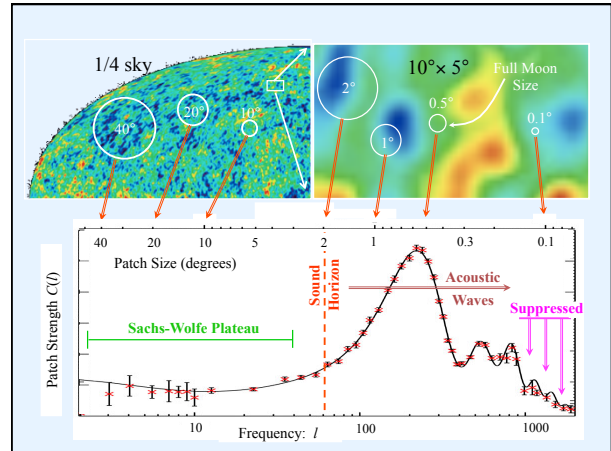
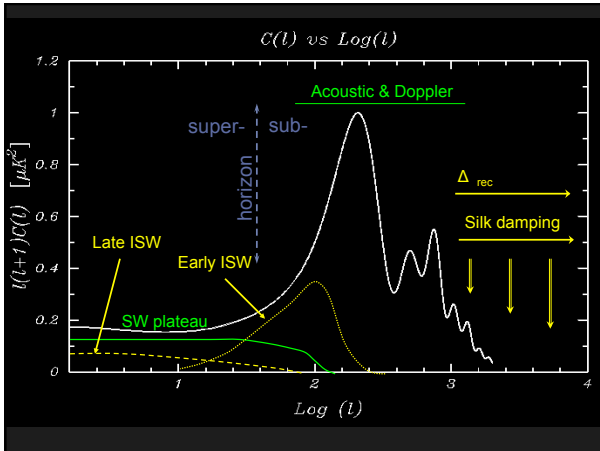
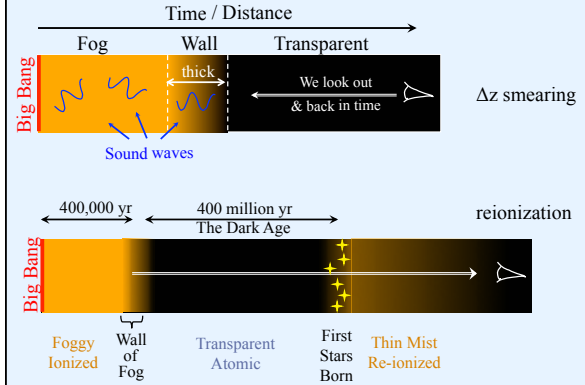
### Origin of Fundamental & Harmonics



## Understanding $C(\ell)$

- **Secondary:**
  1. **Smearing:**  $\lambda < \Delta R_{\text{rec}}$  wash out; kills high  $\ell$
  2. **Silk damping:** photon diffusion, kills high  $\ell$
  3. **Integrated Sachs-Wolfe:**  $\gamma$ s cross varying  $\Phi$ 
    - Early ISW :  $\Phi_\gamma$  near  $z_{\text{rec}}$  (adds power @  $\ell \sim 100$ )
    - Late ISW :  $\Lambda \rightarrow \Phi(z < 1)$  (adds power @  $\ell < 10$ )
  4. **Re-ionization:** @  $z \sim 20$  lowers power for  $\lambda < \lambda_{\text{H}}(z_{\text{re-ion}})$
  5. **Cluster SZ effects:** add power @ high  $\ell > 3000$
- **Tertiary** (contamination):
  1. **Galactic:** dust, free-free, synchrotron
  2. **Point sources:** radio galaxies; high-z IR gals
  3. **Dipole** ( $10^{-3} \times T_{\text{cmb}}$ ;  $10^2 \times$  other anisotropies)

## Two influences at high $\ell$



## Diagnostics: Measuring Cosmic Parameters

Several key datasets are sensitive to the cosmological parameters. These include:

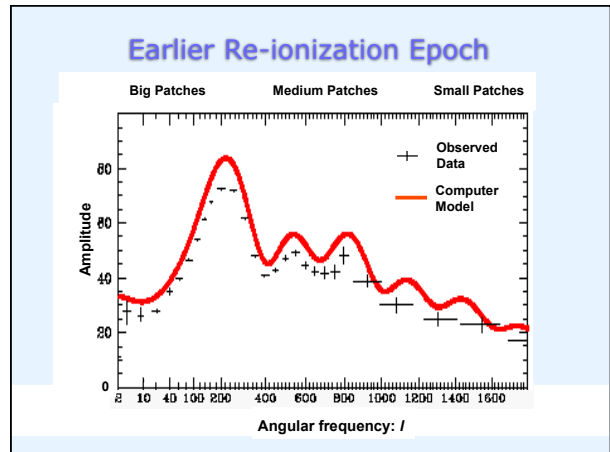
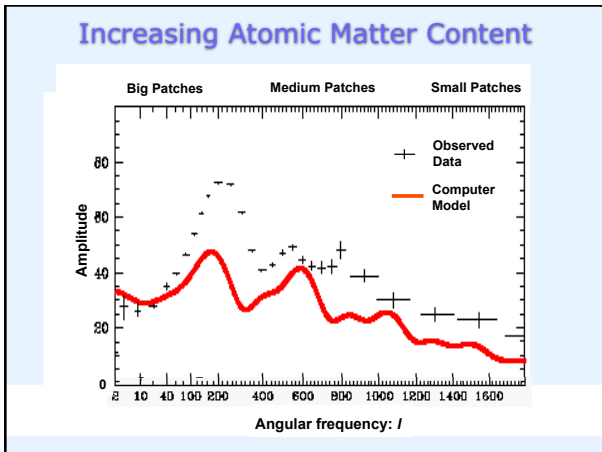
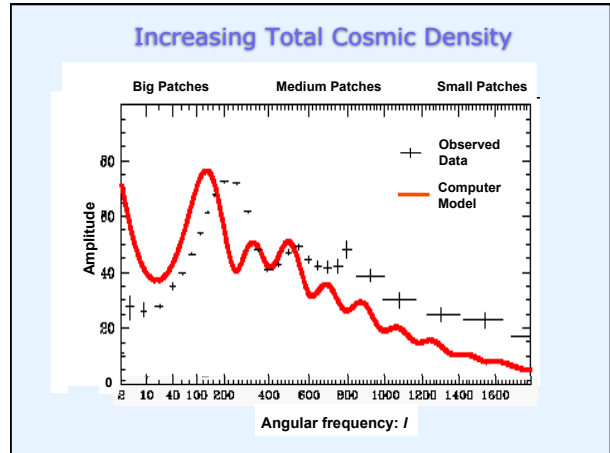
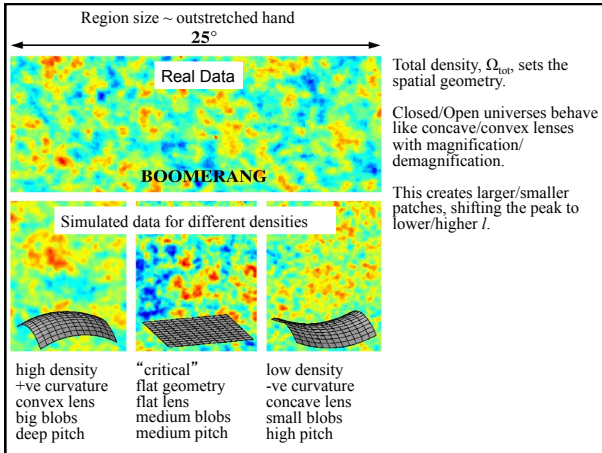
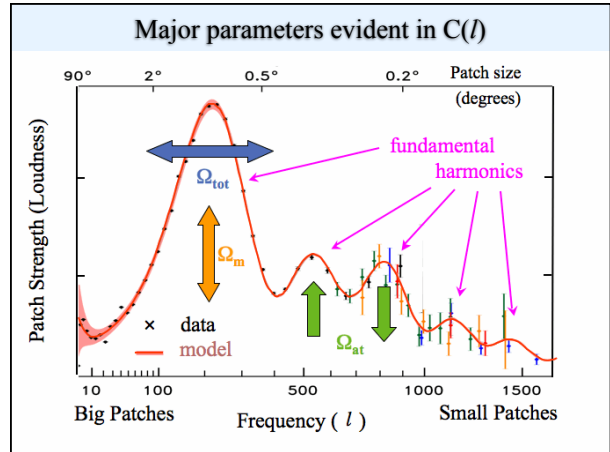
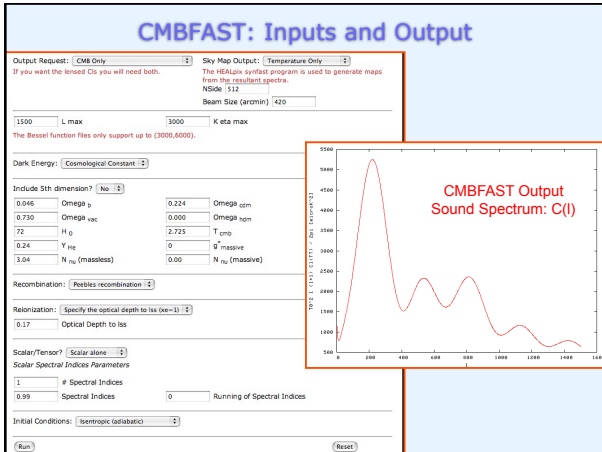
- 1) The CMB temperature, and power spectrum  $C(l)$
- 2) The near-field redshift-distance data, giving  $H_0$ .
- 3) The far-field SNIa redshift-magnitude data.
- 4) The galaxy power spectrum, including BAOs.
- 5) The galaxy cluster number vs redshift relation.
- 6) The abundances of the light elements.

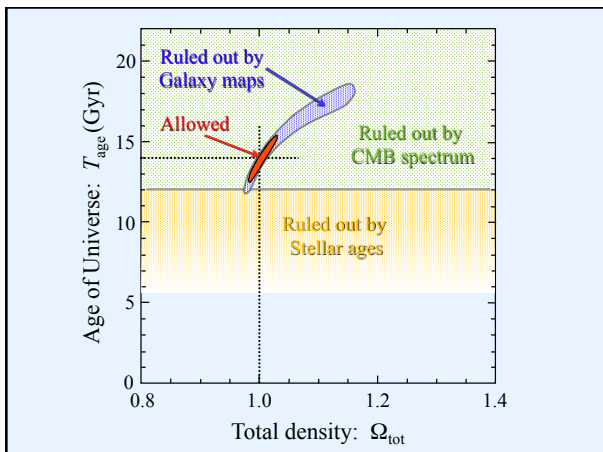
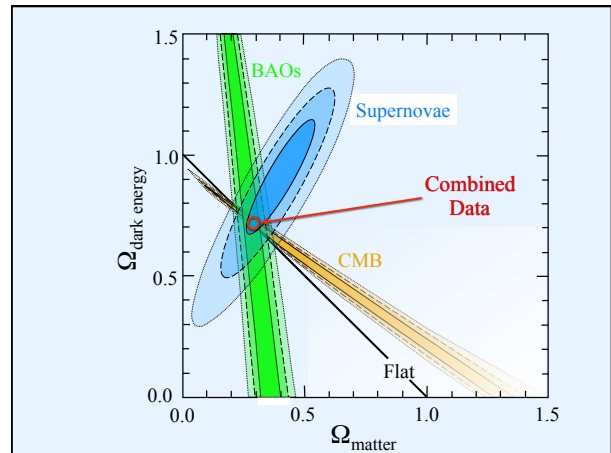
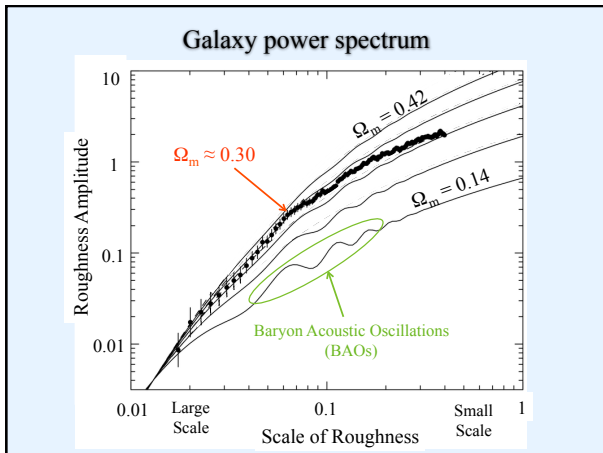
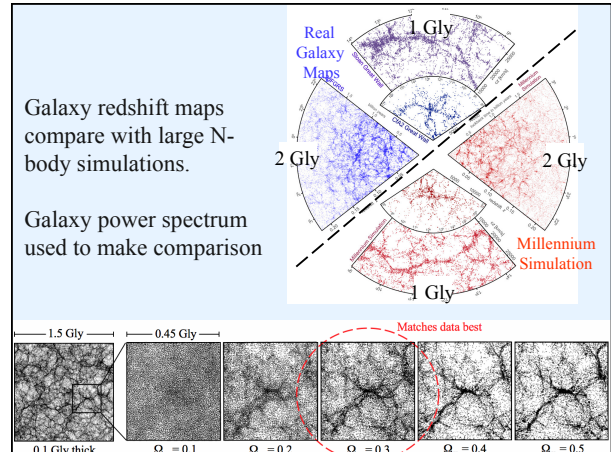
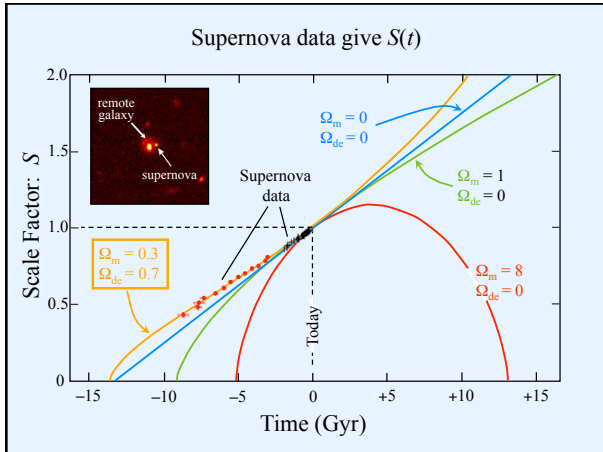
While each dataset primarily measures one or two parameters, they are normally combined to introduce redundancy, and reduce uncertainties in the parameter values.

For some of these – e.g.  $C(l)$ ,  $P(k)$  – one needs a detailed working model to compare to data.

## Modeling $C(\ell)$ and $P(k)$

- **Highly sophisticated; long history:**
  - Early work: Peebles; Silk; Bond; Efstathiou.....
  - Developments: Seljak; Sugiyama; Zaldarriaga; Hu.....
  - Improved numerical methods & physical understanding
  - Public code : CMBFAST (Seljak & Zaldarriaga '96)
- **In theory:**
  - 4 fluids: CDM &  $\nu$ s (collisionless); baryons &  $\gamma$ s (collisional)
  - Evolve fluid DFs using Boltzmann Eqn + perturbations
  - $P(k,z)$  & Transfer functions →  $C(\ell)$
  - Physics “known” → accurate to  $\leq$  few %
- **In practice:**
  - Input global:  $\Omega_b, \Omega_m, \Omega_\Lambda, \Omega_\nu, T_{\text{cmb}}, h$ ; & perturbations:  $n, A, \text{type}$ .
  - Let rip →  $P(k,z)$  for baryons,  $\gamma$ , cdm,  $\nu$ ; +  $C(\ell)$





### Representative parameter values

Best-fit cosmological parameters from WMAP nine-year results<sup>(6)</sup>

Parameter	Symbol	Best fit (WMAP only)	Best fit (WMAP + eCMB + BAO + H <sub>0</sub> )
Age of the universe (Ga)	$t_0$	$13.74 \pm 0.11$	$13.772 \pm 0.059$
Hubble's constant ( $^{km} / Mpc s$ )	$H_0$	$70.0 \pm 2.2$	$69.32 \pm 0.80$
Baryon density	$\Omega_b$	$0.0463 \pm 0.0024$	$0.04628 \pm 0.00093$
Physical baryon density	$\Omega_b h^2$	$0.02264 \pm 0.00050$	$0.02223 \pm 0.00033$
Cold Dark matter density	$\Omega_c$	$0.233 \pm 0.023$	$0.2402^{+0.0089}_{-0.0087}$
Physical cold dark matter density	$\Omega_c h^2$	$0.1138 \pm 0.0045$	$0.1153 \pm 0.0019$
Dark energy density	$\Omega_\Lambda$	$0.721 \pm 0.025$	$0.7135^{+0.0095}_{-0.0096}$
Density fluctuations at $8h^{-1} Mpc$	$\sigma_8$	$0.821 \pm 0.023$	$0.820^{+0.013}_{-0.014}$
Scalar spectral index	$n_s$	$0.972 \pm 0.013$	$0.9608 \pm 0.0080$
Reionization optical depth	$\tau$	$0.089 \pm 0.014$	$0.081 \pm 0.012$
Curvature	$1 - \Omega_{tot}$	$-0.037^{+0.044}_{-0.042}$	$-0.0027^{+0.0039}_{-0.0038}$
Tensor-to-scalar ratio ( $k_B = 0.002 Mpc^{-1}$ )	$r$	$< 0.38$ (95% CL)	$< 0.13$ (95% CL)
Running scalar spectral index	$dn_s/dlnk$	$-0.019 \pm 0.025$	$-0.023 \pm 0.011$